



## RESEARCH ARTICLE

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## Key Points:

- *Trichodesmium* abundance was elevated in certain cyclonic and anticyclonic eddies
- Enhancement in cyclonic eddies could be driven by Ekman convergence
- Anticyclonic eddies with elevated abundance were anomalously fresh

## Supporting Information:

- Supporting Information S1

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Mesoscale eddies and *Trichodesmium* spp. distributions in the southwestern North Atlantic

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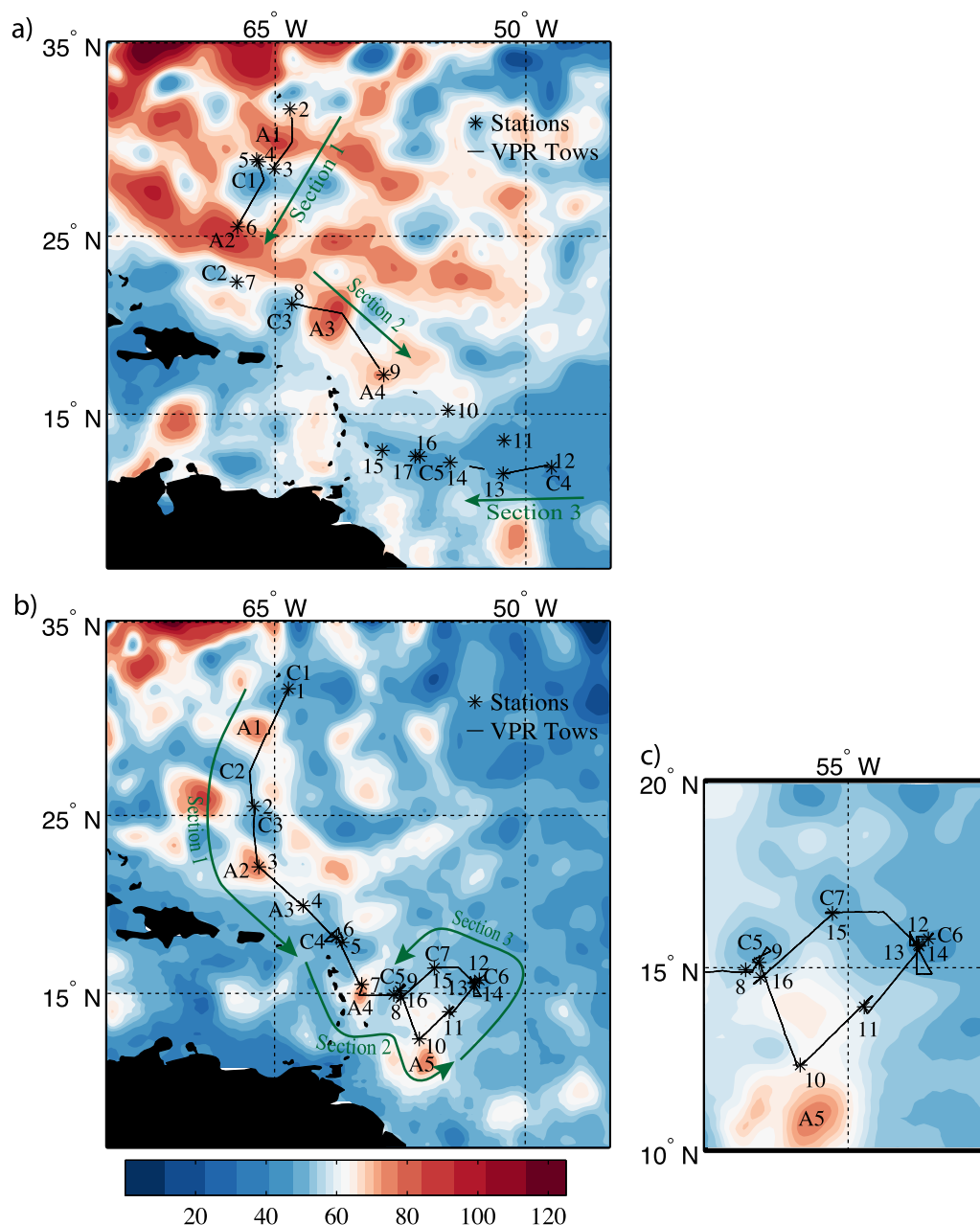
**Abstract** Correlations of *Trichodesmium* colony abundance with the eddy field emerged in two segments of Video Plankton Recorder observations made in the southwestern North Atlantic during fall 2010 and spring 2011. In fall 2010, local maxima in abundance were observed in cyclones. We hypothesized surface Ekman transport convergence as a mechanism for trapping buoyant colonies in cyclones. Idealized models supported the potential of this process to influence the distribution of buoyant colonies over time scales of several months. In spring 2011, the highest vertically integrated colony abundances were observed in anticyclones. These peaks in abundance correlated with anomalously fresh water, suggesting riverine input as a driver of the relationship. These contrasting results in cyclones and anticyclones highlight distinct mechanisms by which mesoscale eddies can influence the abundance and distribution of *Trichodesmium* populations of the southwestern North Atlantic.

## 1. Introduction

*Trichodesmium* is a genus of nitrogen-fixing cyanobacteria found throughout the tropical and subtropical oceans. Converting atmospheric dinitrogen into biologically available forms, it represents an important source of new nitrogen to oligotrophic regions and contributes to primary production [Capone *et al.*, 1997; Carpenter, 1983; LaRoche and Breitbarth, 2005; Letelier and Karl, 1996]. *Trichodesmium* also exports carbon from surface waters [Capone *et al.*, 1998; Dore *et al.*, 2002; Bar-Zeev *et al.*, 2013]. Thus, understanding the controls on *Trichodesmium* abundance and distribution is a step toward quantifying the contribution of *Trichodesmium* to oceanic nitrogen and carbon cycling.

On a Video Plankton Recorder (VPR) survey across the North Atlantic subtropical gyre, Davis and McGillicuddy [2006] identified a correlation between *Trichodesmium* colony abundance and sea level anomaly. Greater abundances were observed in warm and salty anticyclonic eddies. Other observations of elevated *Trichodesmium* abundance in anticyclones have been made in the Pacific [Fong *et al.*, 2008] and eastern North Atlantic [Taboada *et al.*, 2010]. To investigate the spatial and temporal generality of this pattern, we analyzed relationships between *Trichodesmium* colony distributions and the mesoscale eddy field based on sea level anomaly and VPR data from two cruises, in fall 2010 and spring 2011. The cruises took place in the southwestern North Atlantic, typically a region of high *Trichodesmium* abundance.

There was no single, ubiquitous relationship between mesoscale dynamics and colony abundance, and the relationships observed were different from those encountered previously by Davis and McGillicuddy [2006]. However, patterns emerged within localized regions on the two cruises, with elevated abundances in cyclonic eddies during the fall, and in fresh, anticyclonic eddies during the spring. Herein, we examine potential physical and biogeochemical drivers of the two relationships using observations and numerical simulations. A mechanism for population increase through convergent transport of colonies is proposed in cyclonic eddies, while stimulation of population growth through nutrient supply is proposed to explain elevated abundances in fresh, anticyclonic eddies. In concert, this study thus provides a framework for interpreting the drivers of *Trichodesmium* colony distribution in this system. We first present the two types of observed patterns and their proposed drivers separately, and then we examine potential for interaction between the mechanisms and the conditions under which each might dominate.



**Figure 1.** Locations of stations and VPR tows for (a) the Fall Cruise and (b) the Spring Cruise with absolute dynamic topography (ADT) in cm. Black lines indicate VPR tows. Stars indicate station locations. Targeted cyclonic and anticyclonic eddies are labeled with prefixes “C” and “A,” respectively. Labeled sections (arrows) are referred to in subsequent analysis. ADT was calculated by objective analysis of available altimetry data from Envisat and Jason-1/2. (c) Southwestern portion of the Spring Cruise track and altimetry magnified and for a later date (8 May 2011), when cyclone C7 was more apparent.

## 2. Observational Methods

Sampling occurred on two cruises on the R/V *Oceanus*, OC469 from 1–22 October 2010, and OC471 from 23 April to 13 May 2011, beginning near Bermuda and ending near Barbados (Figure 1). Precise sampling locations were chosen during the cruises as informed by satellite observations of sea surface height, real-time analysis of VPR *Trichodesmium* abundance observations, and ocean-color-based estimates of *Trichodesmium* bloom probability [Westberry and Siegel, 2006; Westberry et al., 2005]. This adaptive strategy allowed sampling to be directed toward interesting physical and biological features such as eddies and areas of

predicted *Trichodesmium* abundance. *Trichodesmium* colony abundance was sampled by the VPR, towed between stations (Figure 1). As a result of the timing of station work, a majority of VPR sampling occurred in late afternoon and throughout the night.

The VPR consists of a towed body, containing a CTD and a synchronized video camera and xenon strobe [Davis *et al.*, 2005]. Data are sent in real time through a fiber-optic cable to a shipboard computer system. The VPR undulated between the surface (5–10 m) and approximately 120 m depth, completing an up-down cycle approximately every 6 min, corresponding to one vertical cycle every 1.8 km at 10 knots ( $5.1 \text{ m s}^{-1}$ ). At a frame rate of 30 Hz, each  $984 \times 1009$  pixel video frame was passed through object-identification software to pick out “regions of interest,” which were saved, using a time-stamp naming convention, to a hard disk.

Images were initially sorted using image recognition software [Hu and Davis, 2006]. For the spring cruise, due to the large number of images collected, the sample size was reduced to 1000 images per sampling hour, evenly spaced among the collected images [Olson, 2014]. For both cruises, each machine-classified image was manually checked and reclassified as necessary, into the three *Trichodesmium* morphological categories, puff, raft, and bowtie, and an “other” category containing everything else. *Trichodesmium* puff, raft, and bowtie colonies are readily sampled using the VPR system due to their macroscopic size, but free-living individual trichomes are not distinguishable from other objects of similar shape and size.

The field of view spanned approximately 13 mm  $\times$  13 mm during the fall cruise, whereas during the spring cruise the camera and strobe were oriented to resolve a deeper and narrower field spanning approximately 12 mm by 12 mm in the vertical and horizontal. The corresponding volumes, calibrated before each cruise by the tethered copepod method described by Davis *et al.* [2005], were 6.929 mL and 15.168 mL. These volumes were used to convert observation frequencies to abundances. Abundance data were gridded to bins 5 m deep and 11 and 33 km along the ship track for the fall and spring cruises, respectively. Within each bin, abundances were estimated under the assumption of a Poisson distribution, based on the summed image volume and the number of positive identifications.

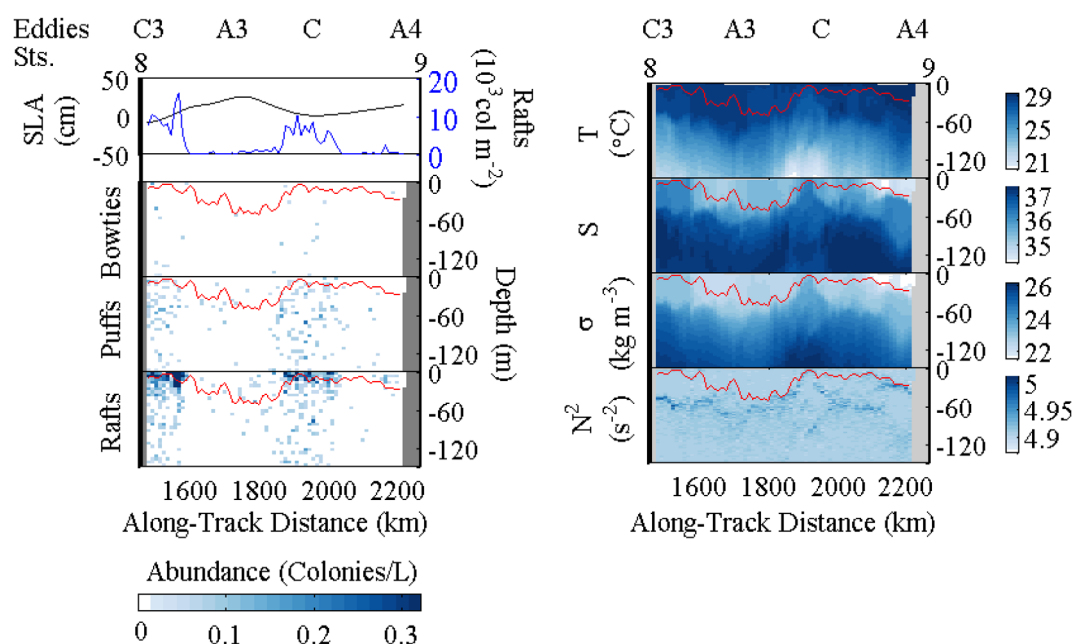
In addition to imaging, the VPR contained a CTD probe that recorded pressure, temperature, salinity, photosynthetically active radiation (PAR), oxygen (on the spring cruise only), and fluorescence. Physical and bio-optical data from the VPR were binned to the same horizontal resolution as the abundance data and either 1 or 5 m vertically, as indicated.

At each hydrographic station in between VPR tows, CTD casts measured temperature, salinity and PAR. 125 ml water samples collected at depths of 700, 500, 300, 200, 100, 80, 60, 40, 20 m, and the surface were filtered through  $0.2 \mu\text{m}$ , 47 mm polycarbonate filters and stored frozen at  $-20^\circ\text{C}$  in 10% HCl-cleaned bottles prior to nutrient analysis. Phosphate [PO<sub>4</sub>] and total dissolved phosphorus [TDP] were analyzed at the Chesapeake Bay Laboratory at the University of Maryland using standard autoanalyzer protocols (nasl.cbl.umces.edu). For samples in the upper 100 m, low level phosphate concentration was measured using a modified MAGIC method [Rimmelin and Moutin, 2005]. Dissolved organic phosphorus (DOP) was assessed as the difference between total and inorganic pools. Alkaline phosphatase activity (APA) was measured using a fluorescent substrate analog after Dyhrman and Ruttenberg [2006]. Measurement detection and quantification limits were: PO<sub>4</sub> DL =  $0.0025 \mu\text{M}$ , QL =  $0.08 \mu\text{M}$ ; TDP DL =  $0.05 \mu\text{M}$ , QL =  $1.45 \mu\text{M}$ .

In the upper 80 m, water samples were gravity filtered and preserved for microscopic enumeration of both *Trichodesmium* colonies and free trichomes. Microscopic colony counts were compared to nearby VPR colony abundances and demonstrated a positive relationship (supporting information Figure S1) [Olson, 2014].

Throughout the cruise, underway sensors measured near-surface environmental conditions. Currents were measured by ADCP. Wind speeds were measured by sensors mounted at 15 m above sea level on the port and starboard sides of the ship. Temperature profiles were recorded using XBTs deployed approximately every hour while underway.

Sea level anomaly (SLA) and absolute dynamic topography (ADT) along the cruise tracks were estimated by linear interpolation in space and time from objectively analyzed fields. The objective analysis combines altimetric data from Envisat, Jason-1 and Jason-2. Eddies were tracked in time and space by visually comparing daily objectively analyzed SLA fields.



**Figure 2.** *Trichodesmium* abundance and physical properties associated with eddies transected on Section 2 of fall 2010 cruise OC469. Black line in top plot is sea level anomaly (cm). Blue line in top plot represents integrated raft colony abundance ( $10^3$  colonies  $m^{-2}$ ). Red lines show mixed layer depth as determined by a threshold criterion of  $0.125 \text{ kg m}^{-3}$  density difference from the surface. Right plots (top to bottom) display temperature, salinity, density, and buoyancy frequency.

### 3. Enhanced Abundance of *Trichodesmium* Colonies in Cyclones

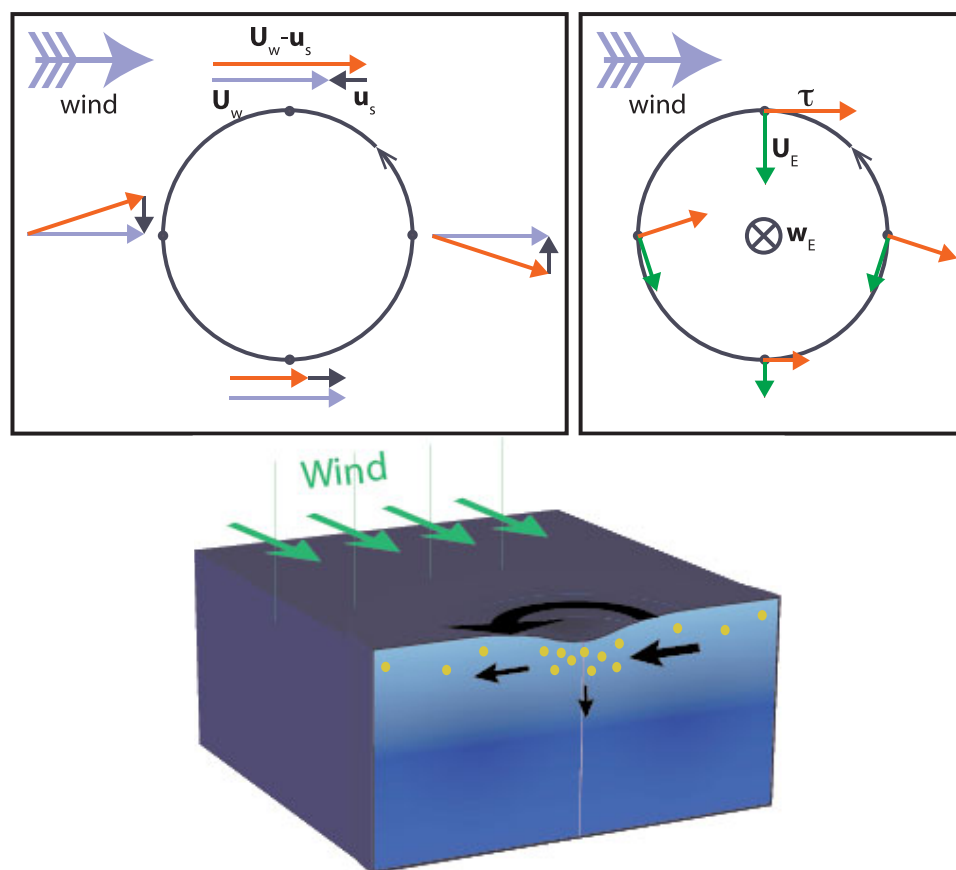
A roughly 700 km region (Section 2 in Figure 1) where *Trichodesmium* colony abundance patterns reflected variations in sea level anomaly associated with the eddy field was transected during the fall cruise. There, in contrast to observations made in Spring 2011 and previously in the subtropical North Atlantic [Davis and McGillicuddy, 2006], *Trichodesmium* colony abundances were locally elevated in cyclones and nearly absent outside of them. Over this region (Section 2, Figures 1 and 2), the correlation coefficient between interpolated sea level anomaly and integrated *Trichodesmium* abundance was  $-0.67$  ( $p < 0.01$ ). Surface abundances within the two cyclones (C3 and C in Figures 1a and 2) fell in the range of  $0.2$ – $0.7$  colonies  $L^{-1}$ , compared to a range from below the limit of detection to  $0.2$  colonies  $L^{-1}$  in the anticyclones (A3 and A4).

Abundances throughout this region were low compared to elsewhere on the cruise, with surface colony abundances averaging  $0.8$  colonies  $L^{-1}$  to the north in Section 1 and  $2.0$  colonies  $L^{-1}$  to the south in Section 3. Rafts were the dominant morphology and were vertically distributed differently from the other morphologies, with greatest abundance within the surface mixed layer. On the southeast side of Eddy C3, depth-integrated colony abundances in the closest four 11 km grid cells to the center averaged  $12.7 \times 10^3$  colonies  $m^{-2}$ . Integrated colony abundances in the seven 11 km grid cells closest to the center of eddy A3 averaged  $0.85 \times 10^3$  colonies  $m^{-2}$ .

At the time of sampling, Eddy C3 had a diameter of roughly 150 km based on the largest closed encircling SLA contour. Although the eddy was not symmetric in shape, a representative radius at which maximum velocities were found was approximately 50 km. Based on shipboard ADCP measurements, maximum velocities at approximately 20–30 m depth were  $0.13 \text{ m s}^{-1}$  on the western side of the eddy and  $0.45 \text{ m s}^{-1}$  on the eastern side.

#### 3.1. Hypothesis: Concentration of Buoyant Particles Through Convergent Ekman Transport in Cyclones

Association of *Trichodesmium* abundance with cyclones was unexpected, as previously greater abundances were observed in anticyclones [Davis and McGillicuddy, 2006; Fong et al., 2008; Taboada et al., 2010]. In some cases, nutrient injection as isopycnals dome upward during eddy formation, with associated shoaling of the nutricline, has been invoked to explain elevated phytoplankton abundance in cyclones



**Figure 3.** Schematic diagram of Ekman convergence concentration of buoyant colonies in a cyclone with wind stress a function of the difference between air and sea surface velocity ( $\tau \propto |U_w - u_s|(U_w - u_s)$ ). (top left) plan view of a cyclonic eddy with the difference (orange) between wind velocity,  $U_w$  (blue), and sea surface velocity,  $u_s$  (black), at four points. (top right) direction of the resulting wind stress,  $\tau$  (orange), and Ekman transport,  $U_E$  (green), at the same locations. Convergence of Ekman transport near the center of the eddy leads to downwelling,  $w_E$ . (bottom) Buoyant colonies become concentrated in a region of surface convergence and downwelling in the center of a cyclonic eddy with uniform wind.

[e.g., Falkowski et al., 1991; McGillicuddy and Robinson, 1997]. Here, however, the inorganic phosphorus pool is deep compared to the location of *Trichodesmium* colonies, found mostly within the mixed layer, and iron is thought to be supplied primarily from surface deposition. In fact, rafts, the most buoyant of *Trichodesmium* colony morphologies, dominated in this region. The buoyancy of *Trichodesmium* colonies and their surface-intensified depth distribution led to a hypothesis of physical concentration by convergent wind-driven surface currents in cyclones.

The hypothesis of wind-driven accumulation of buoyant *Trichodesmium* colonies in cyclones is based on convergent transport in cyclonic eddies due to variations in wind stress across the breadth of an eddy. These wind stress variations result from the difference in the surface currents on opposite sides of the eddy (Figure 3). Since wind stress depends on the relative velocity between air and ocean, it will be greater where the current flows opposite to the wind. As a result, in the surface boundary layer of a cyclone, more water flows into the cyclone on the side where current and wind are in opposition than flows out on the side where current and wind are parallel. This results in surface convergence, balanced by downward motions in the eddy interior. The opposite takes place in anticyclones, with wind-driven surface divergence leading to upwelling. We will refer to this effect as “eddy-wind interaction” throughout. Vertical velocities due to this type of spatial variation in Ekman transport have been identified previously [McGillicuddy et al., 2007; Ledwell et al., 2008] as a mechanism for nutrient supply in anticyclonic and mode-water eddies, and the physical mechanism has been well described [Dewar and Flierl, 1987; Martin and Richards, 2001].



### 3.2. Theoretical Model

Martin and Richards [2001] investigated eddy-wind interaction in the context of an eddy with Gaussian velocity structure at the surface. This idealized eddy surface velocity field reflects solid body rotation at the eddy center and exponential decay beyond the radius of maximum velocity. The azimuthal velocity is

$$\mathbf{v}_\theta(r) = V \frac{r}{R} \exp \left[ \frac{1}{2} \left( 1 - \frac{r^2}{R^2} \right) \right], \quad (1)$$

where  $V$  is the maximum azimuthal velocity, which occurs at a radius  $r = R$  from the eddy center. In the northern hemisphere,  $V$  is positive in a cyclonic eddy.

Using this velocity structure, we calculate the rate of transport of buoyant colonies due to eddy-wind interaction in an idealized system. We assume that the convergent Ekman transport takes place over a uniform mixed layer with *Trichodesmium* concentration  $C$ . The formula for the steady Ekman transport, the depth integral of the ageostrophic velocity induced by the surface stress, is

$$\mathbf{U}_E = (U_E, V_E) = \left( \frac{\tau^{(y)}}{\rho_0 f}, -\frac{\tau^{(x)}}{\rho_0 f} \right) \quad (2)$$

[see Batchelor, 1967, equation (4.4.11)], where  $\rho_0$  is the density of seawater ( $1026 \text{ kg m}^{-3}$ ), and  $f$  is the Coriolis parameter ( $5.21 \times 10^{-5} \text{ s}^{-1}$ ).  $\tau^{(x)}$  and  $\tau^{(y)}$  are the zonal and meridional components of the wind stress at the ocean surface.

Dewar and Flierl [1987] used the wind stress parameterization,

$$\tau = \rho_a c_d |\mathbf{U}_a - \mathbf{u}| (\mathbf{U}_a - \mathbf{u}), \quad (3)$$

where  $\rho_a$  is the density of air ( $1.2 \text{ kg m}^{-3}$ );  $c_d$  is the drag coefficient (0.00118);  $\mathbf{U}_a$  is the wind velocity; and  $\mathbf{u}$  is the sea surface velocity. They assumed a uniform westerly wind and small eddy surface velocities compared to wind velocity ( $|\mathbf{u}| \ll |\mathbf{U}_a|$ ), so that the wind stress expression was simplified to

$$(\tau^{(x)}, \tau^{(y)}) = (\rho_a c_d (U_a^2 - 2U_a u_0), -\rho_a c_d U_a v_0), \quad (4)$$

where  $u_0$  and  $v_0$  are the zonal and meridional components of the azimuthal sea surface velocity associated with the eddy ( $\mathbf{v}_\theta$ , equation (1)). Then, the zonal and meridional components of the Ekman transport are

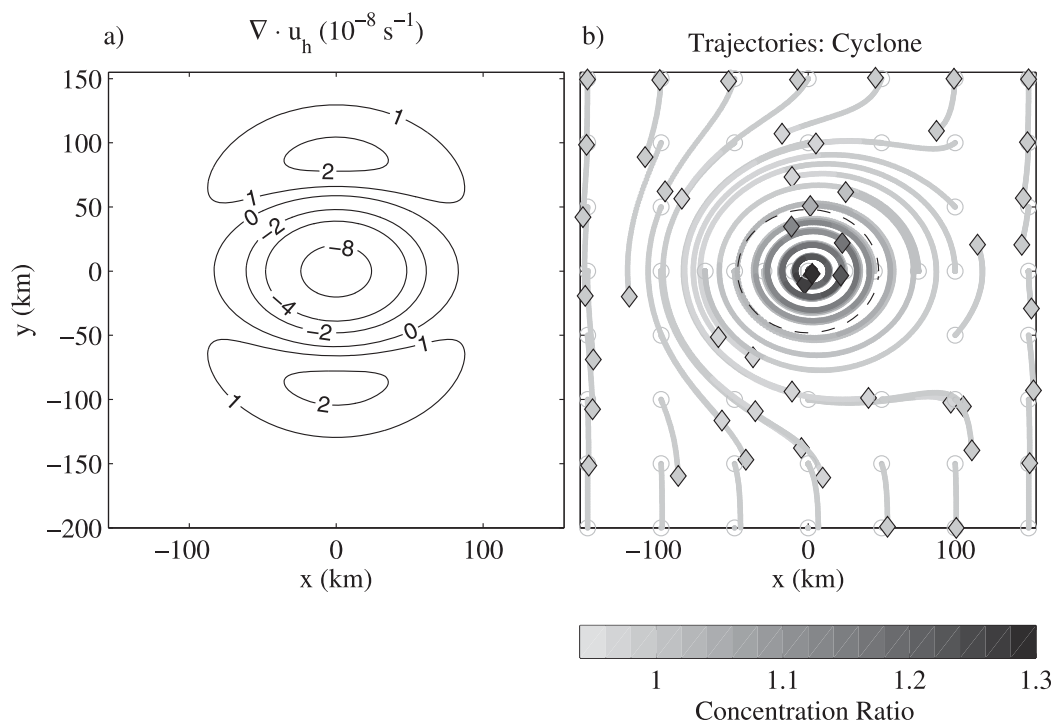
$$U_E = -\frac{\rho_a c_d}{\rho_0 f} U_a v_0, \quad (5)$$

and

$$V_E = -\frac{\rho_a c_d}{\rho_0 f} (U_a^2 - 2U_a u_0). \quad (6)$$

The first term contributing to  $V_E$  (equation (6)) is the nondivergent wind stress in the absence of sea surface motion. The second term in  $V_E$  and the single term contributing to  $U_E$  (equation (5)) reflect modification of the wind stress by the sea surface velocity. In a cyclonic eddy, these terms produce convergence at the eddy center, leading to downwelling there [see Dewar and Flierl, 1987].

Buoyant particles tend to concentrate in regions of surface convergence; they are advected laterally, but float rather than following water trajectories leaving the surface (Figure 3). Similarly, buoyant particles become spread out in regions of surface divergence. This phenomenon has previously been investigated in the context of frontal dynamics [Franks, 1992]. Additionally, submesoscale modeling of effects of surface convergence and divergence on buoyant particle fields has demonstrated concentration in regions of convergence in the presence of an eddying field [Zhong et al., 2012]. However, the influence of sea surface velocity on the wind stress was not represented in the Zhong et al. [2012] model, and therefore it cannot simulate the eddy-wind interaction investigated here. With eddy-wind interaction, we expect the convergence of Ekman transport to lead to concentration of *Trichodesmium* colonies in cyclones (Figure 3), as previous investigation of *Trichodesmium* colony floating velocities indicated that a majority of colonies were positively buoyant [Walsby, 1978]. Thus, the effect could lead to higher concentrations of buoyant *Trichodesmium* colonies in cyclones than anticyclones, consistent with the observed distribution of *Trichodesmium* colonies in relation to the mesoscale eddy field in fall 2010.



**Figure 4.** (a) Contours of the divergence of the horizontal velocity field (equations (8) and (9)). Flow is convergent near the eddy center, with regions of divergence to the north and south in the presence of a uniform westerly wind. (b) 30 day trajectories calculated based on the same horizontal surface boundary layer velocities, which result from a combination of azimuthal eddy velocities and wind-driven Ekman transport. A black dashed line marks the radius of maximum velocity of the eddy at 48 km. A circle represents the point of origin of each trajectory, and a diamond marks the end location. Shading indicates tracer concentration as the ratio of final to initial value along each Lagrangian trajectory. Tracer concentration was calculated by integration of equation (7), which does not include the effects of down-gradient diffusive flux due to horizontal mixing.

We here present an idealized analytical model for the change in concentration of buoyant tracer over time, based on the convergence of Ekman transport near the center of a cyclonic eddy (section 3.1). For the purpose of these calculations, we supposed that the horizontal Ekman transport occurs over an initially vertically homogeneous surface layer of depth  $H$  containing uniform tracer concentration. We further supposed that over time, the tracer maintains its position vertically, and is thus subject to horizontal but not vertical advection. Thus, in the absence of diffusive mixing, the fractional rate of change in time of concentration along a Lagrangian path is proportional to the opposite of the divergence of the velocity field.

$$\frac{1}{C} \frac{DC}{Dt} = -\nabla \cdot \mathbf{u}_h, \quad (7)$$

where  $\nabla \cdot \mathbf{u}_h$  is the divergence of the horizontal velocity field,  $C$  is concentration, and  $\frac{D}{Dt}$  denotes the Lagrangian (or material) derivative. In regions of convergence, tracer maintains its position as the fluid stretches vertically, so tracer concentration increases. In regions of divergence, tracer-free water is injected into the upper layer, decreasing local tracer concentration. The velocity in this layer,  $\mathbf{u}$ , with components  $u$  and  $v$ , is the sum of the azimuthal eddy surface velocity (equation (1)) and the Ekman transport (equations (5) and (6)) divided by the layer depth,  $H$ :

$$u = u_0 + u_E = -v_0 \sin \theta - \frac{\rho_a C_d}{\rho_0 f H} U_a v_0 \quad (8)$$

$$v = v_0 + v_E = v_0 \cos \theta - \frac{\rho_a C_d}{\rho_0 f H} (U_a^2 - 2U_a u_0) \quad (9)$$

Integration of this velocity field reveals that in a cyclone, trajectories originating close to the eddy center are trapped within the eddy, slowly spiraling inward toward the center (Figure 4).

**Table 1.** Constants and Eddy Characteristics

Description	Variable	Value	Units
Density, air	$\rho_a$	1.2	$\text{kg m}^{-3}$
Density, seawater	$\rho_0$	1026	$\text{kg m}^{-3}$
Coriolis parameter, 21° N	$f$	$5.21 \times 10^{-5}$	$\text{s}^{-1}$
Drag coefficient	$c_d$	$1.18 \times 10^{-3}$	
Surface wind speed	$U_a$	4.7	$\text{m s}^{-1}$
Eddy radius of maximum velocity	$R$	48	km
Maximum azimuthal velocity, cyclone	$V$	0.222	$\text{m s}^{-1}$
Mixed layer depth	$H$	30	m

We were interested in the overall change in tracer concentration over the eddy core, which we defined as the region within the radius of maximum velocity ( $r < R$ ). The horizontal velocity divergence is equal to the Ekman pumping,  $w_e$ ,

$$w_e = \frac{\rho_a}{\rho_0 f_0} c_d U_a (2(u_0)_y - (v_0)_x) \quad (10)$$

[Dewar and Flierl, 1987].

We estimated the horizontal flux of water into an eddy,  $F$ , as the opposite of the Ekman pumping integrated over the eddy core.

$$F = \frac{\rho_a}{\rho_0 f_0} c_d U_a 3\pi R V. \quad (11)$$

At time  $t = 0$ , if *Trichodesmium* colonies were spread uniformly throughout the surface layer with a concentration  $C_0$ , the initial flux of colonies into the eddy center would be  $\phi = C_0 F$ . Then, the fractional rate of change of integrated abundance within the eddy would be

$$\frac{1}{C_0} \frac{dC}{dt} = \frac{3\rho_a c_d U_a V}{\rho_0 f_0 R H}. \quad (12)$$

Using parameters from Table 1, based on Eddy C3, the initial relative rate of increase of concentration within the center of the cyclone would be 0.5% per day.

Over time, additional factors influence the flux of tracer into the eddy. The tracer concentration itself develops spatial variability, which precludes application of equation (12). As gradients in tracer concentration develop, diffusive processes will tend to counteract them. Additionally, on the northern and southern flanks of the eddy (Figure 4a), the horizontal velocity field described by equations (8) and (9) has regions of divergence. Tracer concentrations decrease along trajectories as they pass through those regions, eventually influencing conditions in the eddy as those trajectories continue into the center region.

Many factors not represented in this simplified analytical model could be simulated more realistically in a numerical model. These include time dependent eddy dynamics (e.g., eddy decay) as well as the effects of continuous and evolving spatial variability in tracer distribution. A number of simplifying assumptions were made in this derivation, including constant mixed layer depth, uniform vertical distribution of colonies within the mixed layer, and Ekman transport wholly within the mixed layer, that are relaxed in the numerical simulations described next. Additionally, small differences were anticipated due to the neglected higher order terms in the wind stress formulation (equations (3) and (4)).

### 3.3. Numerical Model

The rate of wind-driven accumulation of buoyant particles, such as *Trichodesmium* colonies, in the center of a cyclonic eddy was investigated with a series of numerical simulations using a three-dimensional primitive equation model. The core numerical experiment consisted of a pair of simulations of cyclonic eddies with different surface wind stress formulations. In one simulation, surface stress was calculated based solely on surface wind velocities. In the other, more realistic, parameterization, surface stress was related to the relative air-sea velocity at the surface (equation (3)). Thus, in the first simulation, the mechanism illustrated in Figure 3 and referred to throughout as “eddy-wind interaction” was turned off, and in the second it was turned on. Buoyant colonies were represented by a tracer with an imposed upward velocity. By comparing the two simulations, we assessed the extent to which spatial variation in wind stress due to eddy surface velocities could drive concentration of buoyant colonies in cyclonic eddies. Additional simulations were carried out to investigate the sensitivity of the results to the floating velocity and strength of the simulated eddy as well as the behavior of the eddy in the absence of wind (Table 2).



**Table 2.** Summary of Model Runs

Label	Eddy-Wind Interaction	Tracer Floating Velocity	Eddy Strength Factor	Wind Speed
$\tau(U_w)$	off	$1.2 \times 10^{-4} \text{ m s}^{-1}$ (10 m d <sup>-1</sup> )	1	$4.7 \text{ m s}^{-1}$
$\tau(U_w - U_s)$	on	$1.2 \times 10^{-4} \text{ m s}^{-1}$ (10 m d <sup>-1</sup> )	1	$4.7 \text{ m s}^{-1}$
w1	on	$1.2 \times 10^{-5} \text{ m s}^{-1}$ (1 m d <sup>-1</sup> )	1	$4.7 \text{ m s}^{-1}$
w100	on	$1.2 \times 10^{-3} \text{ m s}^{-1}$ (100 m d <sup>-1</sup> )	1	$4.7 \text{ m s}^{-1}$
strong	on	$1.2 \times 10^{-4} \text{ m s}^{-1}$ (10 m d <sup>-1</sup> )	1.6	$4.7 \text{ m s}^{-1}$
no wind		$1.2 \times 10^{-4} \text{ m s}^{-1}$ (10 m d <sup>-1</sup> )	1	0

equation model employing the hydrostatic and Boussinesq approximations [Smith *et al.*, 2000]. Horizontal mixing of momentum and tracers was Laplacian with diffusivity and kinematic viscosity of  $1 \text{ m}^2 \text{ s}^{-1}$ . The model included k-profile vertical mixing with a background diffusivity of  $10^{-5} \text{ m}^2 \text{ s}^{-1}$ . Tracer advection was by a centered difference scheme (MPDCD) with a flux limiter to prevent negative concentrations. Simulations were carried out on an f-plane, at  $21^\circ \text{ N}$ , to avoid westward translational motion of the eddies due to the variation in planetary vorticity with latitude and facilitate model diagnosis in eddy-centric coordinates. The vertical grid contained 72 levels with resolution increasing from 125 m at 4462 m depth to 5 m at the surface. Horizontally, the doubly periodic domain contained  $141 \times 141$  cells with zonal and meridional resolution of approximately 3.1 km.

We used the POP version 2 free surface boundary condition. Restoring freshwater and heat fluxes at the surface were applied on time scales of 1 day. Wind stress was formulated as in Ledwell *et al.* [2008], with parameterizations for uniform and surface-current-dependent stress. Wind speed was constant at  $4.7 \text{ m s}^{-1}$ , rotating in direction on a 64 h cycle to avoid net Ekman transport in the horizontal direction. The chosen wind speed was the average wind speed observed during the fall cruise for the period 9–15 October 2010, when the ship was near the region where elevated *Trichodesmium* colony abundance was observed in cyclones.

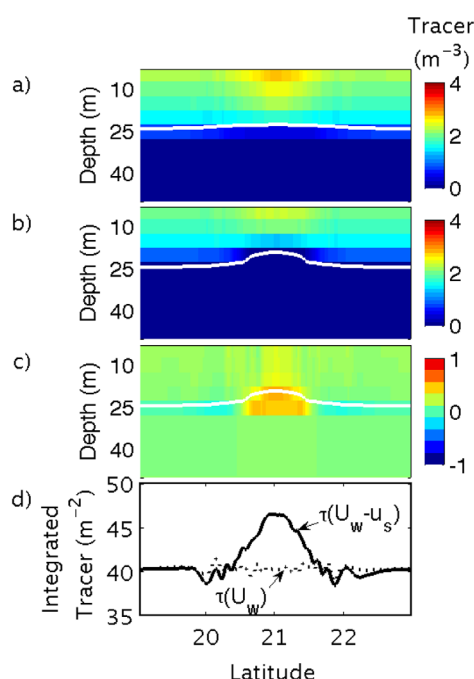
As in Ledwell *et al.* [2008], hydrographic profiles from the center and periphery of an observed eddy were used to initialize the model temperature and salinity fields (supporting information Figure S2), with Gaussian interpolation between center and edge profiles with an e-folding scale of 68 km. The eddy simulated was cyclonic Eddy C3 from the fall cruise (Figure 1a). A station was located at the center of Eddy C3, so CTD measurements were available there to 700 m. CTD temperature and salinity measurements from the VPR were available in the upper 150 m at the edge location. At the edge, XBT temperature measurements were available to 700 m. Salinity between 150 m and 700 m was inferred based on the temperature-salinity relationship observed at Station 8, at the eddy center. Below 700 m at both center and edge, climatological values based on the World Ocean Atlas 2009 [Antonov *et al.*, 2010; Locarnini *et al.*, 2010] were used.

Initial velocity fields were calculated based on geostrophic and cyclostrophic balance with zero motion at the lowermost model level. Whereas addition of a barotropic velocity component brought velocity fields into agreement with observations in the Ledwell *et al.* [2008] eddy simulation, such an addition was not required here. The velocity structure was characterized by solid body rotation in the eddy core transitioning to exponential decay with increasing distance from the eddy center. A maximum azimuthal velocity of  $0.22 \text{ m s}^{-1}$  occurred at a radius of 48 km, roughly consistent with ADCP measurements of velocities associated with Eddy C3.

*Trichodesmium* colony buoyancy was simulated by imposing an upward vertical tracer velocity at every location except the surface boundary. Buoyant *Trichodesmium* colonies have been observed to rise at velocities greater than  $0.001 \text{ m s}^{-1}$  (90 m d<sup>-1</sup>) [Walsby, 1992]. However, a large range of velocities may be measured in a given location, some positive and some negative. In one study [Walsby, 1978], 69% of *T. thiebautii* (puff) colonies, 87% of *T. contortum* (raft) colonies, and 100% of *T. erythraeum* (raft) colonies were positively or neutrally buoyant. In a sample collected at 20 m, Walsby [1978] measured floating velocities of *T. thiebautii* (puff) colonies with a range of  $\pm 7 \times 10^{-4} \text{ m s}^{-1}$  and a mean velocity of  $-3 \times 10^{-5} \text{ m s}^{-1}$ . Mean observed raft velocities were not specified, but it seems reasonable to assume they were positive, based on the greater proportion of floating raft colonies compared to puff colonies. Based on these data, a model tracer velocity of  $1.2 \times 10^{-4} \text{ m s}^{-1}$  (10 m d<sup>-1</sup>) was chosen for the present study, with sensitivity analyses conducted at  $1.2 \times 10^{-3}$  and  $1.2 \times 10^{-5} \text{ m s}^{-1}$ .

### 3.3.1. Model Configuration

Each simulation contained an isolated eddy within a doubly periodic domain. The model implementation was nearly identical to that described by Ledwell *et al.* [2008]. The model was the Los Alamos Parallel Ocean Program (POP) version 2, a three-dimensional primitive



**Figure 5.** 20 day (Days 12–31) mean tracer concentrations. (top three plots) tracer concentration fields in a cross section through the eddy center with white contours marking mixed layer depths determined by a  $0.025 \text{ kg m}^{-3}$  potential density difference from the surface. (a) With eddy-wind interaction:  $\tau(U_w - u_s)$ . (b) Without eddy-wind interaction:  $\tau(U_w)$ . (c) Difference between  $\tau(U_w - u_s)$  and  $\tau(U_w)$ . Mixed layer depths are shown in white; the  $\tau(U_w)$  mixed layer depth is repeated in Figure 5c. (d) Depth-integrated tracer concentration. Solid line:  $\tau(U_w - u_s)$ . Dashed line:  $\tau(U_w)$ .

ing the time for the tracer to adjust to its quasi-equilibrium vertical distribution, a balance between floating velocity, vertical fluid velocity, and vertical mixing. This vertical distribution was achieved within the first day of simulation, so that by day 2 the daily average vertical distribution of colonies was within 1.4% of its quasi-equilibrium distribution, based on the distribution at the end of the initial adjustment period.

### 3.3.2. Numerical Model Results

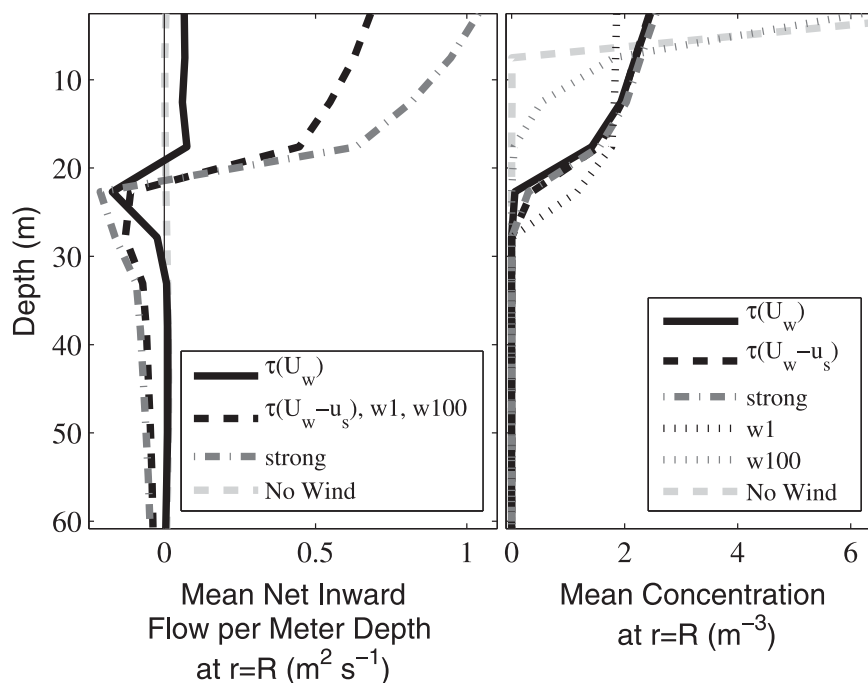
Time-mean cross sections revealed the impact of eddy-wind interaction on simulated tracer distributions (Figure 5). In the simulation without eddy-wind interaction, mixed layer depth was shallower inside the core of the cyclone than outside it. Therefore, buoyant tracer mixed vertically over a shorter distance inside the cyclone, producing a local maximum in surface tracer concentration at the eddy center. This is why concentration increased at the surface in the center of the eddy, even though integrated tracer concentrations were roughly constant throughout the model domain. Net flux into the eddy across a circle of radius of 48 km (approximately the radius associated with maximum eddy velocity) was nearly zero throughout the upper 50 meters (Figure 6).

Compared to the cyclone without eddy-wind interaction, the cyclone with eddy-wind interaction exhibited less doming of the pycnocline as demonstrated by the mixed layer depth (Figure 5). This flattening of isopycnals resulted from downward motion within the eddy due to convergent Ekman transport. The inward volume flux associated with these downward motions was concentrated in the upper 25 m (Figure 6) and led to accumulation of buoyant tracer near the cyclone center. This horizontal convergence of buoyant tracer was associated with a maximum in integrated tracer concentration at the center of the eddy (Figure 5, bottom left). In both cases, with and without eddy-wind interaction, the mean mixed layer depth deepened over the course of the simulations, from initial conditions with weak surface stratification to a mixed layer of approximately 30 m by day 50.

We assessed the accumulation of tracer as the percent difference in mean integrated concentration within the core ( $r < 48 \text{ km}$ ) from that of the initial condition (Figure 7). In the cyclone without eddy-wind

As the velocity fields used to initialize the model were in approximate balance, an initial adjustment of velocity and density structure took place at the beginning of the model run, with the largest changes taking place on the first day. To avoid influence of this initial adjustment period on the results, the tracer concentration was reset after 8 days of simulation and the model restarted. Therefore, in the results that follow, time is referenced to the restart day rather than the day on which the model physics were initialized. As a result, maximum azimuthal eddy velocities on day 1 differ from those calculated in the initial conditions, changing from  $0.22$  to  $0.23 \text{ m s}^{-1}$  in the cyclone with eddy-wind interaction and to  $0.24 \text{ m s}^{-1}$  in the cyclone without eddy-wind interaction.

The tracer initial condition consisted of constant concentration throughout the upper 20 m of the model domain. This placed the tracer initially within the mixed layer, reducing



**Figure 6.** (left) 20 day mean vertical profiles of inward volume transport per meter depth across a circle centered at eddy center with walls at a radius of 48 km, approximately the radius of maximum eddy velocity. (right) 20 day mean vertical profiles of tracer concentration at 48 km from eddy center. Volume transports and concentration profiles are displayed for runs described in Table 2.

interaction, the difference in tracer concentration was negligible throughout the model run. In the cyclone with eddy-wind interaction, the percent difference increased throughout the simulation.

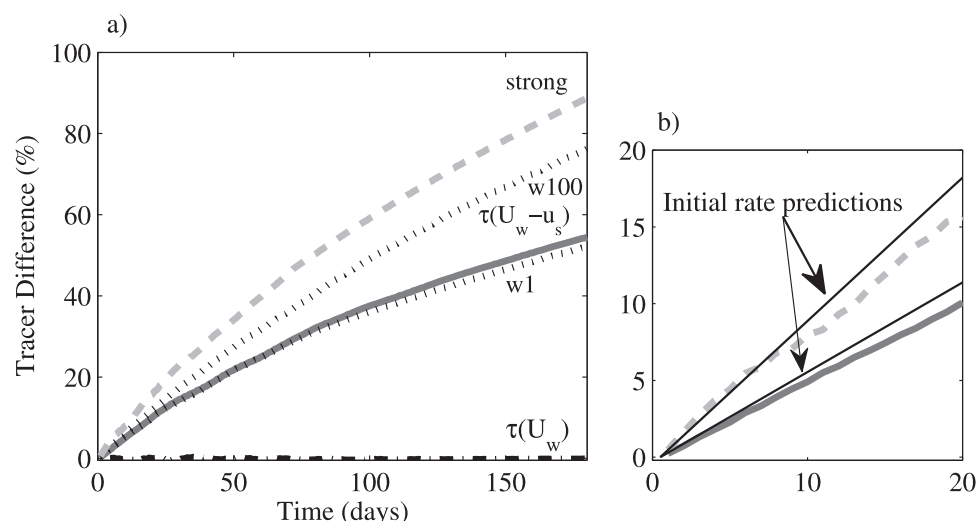
The analytic model presented in section 3.2 provides an explicit prediction for the initial rate of accumulation (equation (12)). For a Gaussian eddy with maximum velocity consistent with day 1 of the numerical simulation ( $0.2316 \text{ m s}^{-1}$ ), the predicted rate was 0.58% per day. The initial rate of concentration in the " $\tau(U_w - u_s)$ " simulation was 0.50% per day (Figure 7b). For the "strong" case, the predicted rate was 0.93% per day and initial rate in the simulation was 1.1% per day. The agreement between the analytic model and the early part of the numerical simulation suggests the underlying conceptual model of accumulation is correct.

The rate of change of tracer concentration in the numerical model slowed over time (Figure 7a) due to a combination of factors. In part, it could be explained by the decaying eddy velocities, which were accompanied by a reduction in the divergent component of the Ekman transport. As discussed by *Dewar and Flierl* [1987], the eddy-wind interaction mechanism ("top drag") contributes to eddy decay. Total kinetic energy decayed over time in all simulations (supporting information Figure S3). However, in the simulations including the eddy-wind interaction mechanism, the rate of decay was more than four times as rapid as in the simulation without eddy-wind interaction.

Another mechanism leading to slowing of the rate of accumulation of tracer was down-gradient diffusive flux, which increases as tracer gradients grow. Accounting for down-gradient diffusive flux of tracer due to unresolved turbulent eddies, the equation governing tracer concentration along a Lagrangian trajectory (section 3.2) becomes

$$\frac{1}{C} \frac{DC}{Dt} = -\nabla \cdot \mathbf{u}_E - \frac{1}{C} A_H \nabla^2 C, \quad (13)$$

where  $A_H$  is a coefficient used to parameterize horizontal mixing. Along a trajectory, changing tracer concentration is controlled by horizontal divergence of the velocity field, producing gradients in tracer concentration, and diffusive processes, smoothing those gradients. The importance of the diffusive flux term ( $\frac{1}{C} A_H \nabla^2 C$ ) was evaluated by comparison with the convergence of the velocity field ( $-\nabla \cdot \mathbf{u}_E$ ). Based on the



**Figure 7.** (a) Difference in eddy center tracer concentration from initial concentration versus time. Center concentration is the average depth-integrated tracer concentration within a radius of 48 km of the eddy center. Lines represent various model runs described in Table 2: " $\tau(U_w)$ " (dot-dashed), " $\tau(U_w - u_s)$ " (solid), "w1" (dotted, bottom), "w100" (dotted, top), and "strong" (dashed). (b) Enlarged lower right corner of Figure 7a, showing only the " $\tau(U_w - u_s)$ " and "strong" cases, which have the eddy-wind interaction turned on and tracer floating velocities of  $10 \text{ m d}^{-1}$ . Thin black lines represent predictions of initial accumulation rates calculated based on Ekman pumping in idealized eddies (equation (12) with eddy velocities of  $v_{\theta} = 0.23 \text{ m s}^{-1}$  ( $\tau(U_w - u_s)$ ) and  $v_{\theta} = 0.37 \text{ m s}^{-1}$  ("strong"), representative of model geostrophic azimuthal velocities on simulation day 1.

analytical velocity field, the maximum contribution by convergence was  $9.9 \times 10^{-8} \text{ s}^{-1}$ . We estimated the diffusive flux based on the difference in tracer between the eddy center and a distance of approximately  $2R$ , or about 100 km. We used a diffusivity of  $A_H = 200 \text{ m}^2 \text{ s}^{-1}$  at a scale of 100 km, based on Okubo [1971]. This was an "effective" diffusivity representing subeddy-scale turbulent transports. The approximate compensating contribution by diffusive processes based on the tracer field in the numerical solution was  $-9.2 \times 10^{-9} \text{ s}^{-1}$  on day 55 of the simulation and  $-1.9 \times 10^{-8} \text{ s}^{-1}$  on day 180. Based on the reduction in maximum geostrophic velocities, the peak convergence would have fallen to  $7.3 \times 10^{-8} \text{ s}^{-1}$  on day 55 and  $3.9 \times 10^{-8} \text{ s}^{-1}$  on day 180. Thus the estimated impacts of diffusion increase over time to nearly 50% of the peak convergence based on the idealized velocity field. Over time, eddy velocities decay, leading to a decrease in the magnitude of the convergence, while the diffusive flux increases with the tracer gradient.

On longer time scales, trajectories crossing into the center of the eddy might eventually carry less tracer if they originated far enough from the eddy center due to dilution of tracer in regions of divergence on the northern and southern edges of the eddy (Figure 4a). However, integrating tracer concentration along trajectories using the analytical model (equations (8), (9), and (13)), in the absence of diffusive effects, suggested this effect would not become evident until after about 140 days, near the end of these simulations.

Under the influence of eddy-wind interaction, tracer concentrations increased by 54% at the center of the cyclonic eddy over the course of the 6 month simulation, for the basic case (" $\tau(U_w - u_s)$ "), compared to a negligible difference in concentration at the end of the run without eddy-wind interaction (" $\tau(U_w)$ ") (Table 3). Additional simulations addressed the sensitivity of this result to uncertain parameters in the model. Analysis of sea level anomaly fields over time suggested that fall Eddy C3 may have been approximately 1.6 times stronger several months prior to sampling. Amplification of the eddy signal by that factor ("strong") increased the tracer concentration effect to 89% over the 6 month simulation. A factor of 10 increase in

floating velocity ("w100") increased the ultimate change in concentration to 76%, while a factor of 10 decrease in floating velocity only decreased the concentration change to 52%.

**Table 3.** Percent Change in Tracer Concentration After 6 Months of Simulation

Label	Tracer Increase
$\tau(U_w - u_s)$	+54%
$\tau(U_w)$	-0.2%
w1	+52%
w100	+76%
strong	+89%

### 3.3.3. Comparison with Observations

In order to compare the amplitude of the increase in the model with observations, we estimated the percent

difference in integrated colony abundance within fall Eddy C3 from background levels (Table 4). The abundance inside C3, estimated based on the four 11 km bins of VPR track closest to the eddy center, was approximately  $1.27 \times 10^4$  colonies  $\text{m}^{-2}$ . In the observations, the assessment of a background level of integrated *Trichodesmium* abundance was not as straightforward. We therefore estimated it in several different ways. First, we computed the average abundance over the segment of VPR track beginning near the center of C3 to the point on the track within A3 with the highest sea level anomaly, as interpolated in time and space from objectively analyzed fields (C3–A3, Table 4). This background estimate, combined with the mean abundance inside C3, corresponds to a percent increase of 153% within the cyclone. This is higher than the 111% increase that might be anticipated in a numerical simulation in which the “strong” case were combined with high floating velocity, under the assumption that the increase in tracer accumulation due to those modifications would add linearly.

However, the C3–A3 background estimate may be skewed low due to the greater track length sampled within A3 compared to C3. Therefore, we also estimated background abundance as the average of the mean integrated abundances within C3 and A3, where the abundance within A3 was calculated based on the seven 11 km VPR track points nearest the eddy center. This corresponds to an increase of 87.6% within the cyclone, comparable to the level of accumulation at the end of 6 months in the “strong” simulation.

We also estimated background abundances based on data compiled by Luo *et al.* [2012] collected between the months of August and December of all years and in the region from  $20^\circ \text{N}$ – $29^\circ \text{N}$  and  $70^\circ \text{W}$ – $59^\circ \text{W}$  (“climatology,” Table 4 and supporting information Figure S4). We chose that region, which is centered to the north of the C3 and A3 sampling site, taking into account the observed patterns in integrated *Trichodesmium* abundance. Abundance in the Luo *et al.* [2012] climatology increased significantly to the south of the C3 and A3 site, a pattern likely attributable to freshwater influence (see section 4), which was minimal in the vicinity of C3 and A3. The percent increase relative to the arithmetic mean of the selected data was 59.2%, and the increase was 107% compared to the geometric mean. The former level is reached by the basic “ $\tau(U_w - u_s)$ ” case, and the latter by the “strong” case.

**Table 4.** Fall 2010 Integrated *Trichodesmium* Colony Abundances<sup>a</sup>

Source	Metric	Value
Inside C3, nearest four 11 km bins	mean	12700 colonies $\text{m}^{-2}$
	standard dev.	1400 colonies $\text{m}^{-2}$
	n	4
Inside A3, nearest seven 11 km bins	mean	852 colonies $\text{m}^{-2}$
	standard dev.	261 colonies $\text{m}^{-2}$
	n	7
Background estimate, observations (C3–A3)	mean	5010 colonies $\text{m}^{-2}$
	standard dev.	5920 colonies $\text{m}^{-2}$
	n	25
Background estimate, climatology	mean	7950 colonies $\text{m}^{-2}$
	standard dev.	6520 colonies $\text{m}^{-2}$
	n	10
	geometric mean	6120 colonies $\text{m}^{-2}$
Percent increase inside C3, relative to:	observations (C3–A3)	153%
	climatology, mean	59.2%
	climatology, geometric mean	107%
	average of C3 and A3 centers	87.4%

<sup>a</sup>Abundance inside C3 was calculated based on the nearest four 11 km VPR data bins to the eddy center. Abundance inside A3 was calculated based on seven 11 km VPR track bins nearest to the eddy center. n is either the number of 11 km bins included in the calculations based on VPR observations, or the total number of measurements for the calculations based on Luo *et al.* [2012]. Background abundance was estimated as the average along a transect beginning inside C3 and ending at the highest sea level anomaly point sampled in A3. Climatological estimates based on Luo *et al.* [2012] include all available integrated *Trichodesmium* abundances sampled between the months of August and December,  $20^\circ \text{N}$  and  $29^\circ \text{N}$ , and  $70^\circ \text{W}$  and  $59^\circ \text{W}$  (see supporting information Figure S4).

### 3.4. Discussion: Eddy-Wind Interaction and Accumulation in Cyclones

Through idealized eddy simulations based on a three-dimensional primitive equation model, we have shown that Ekman transport convergence due to the influence of eddy velocities on the wind stress field can lead to increased buoyant tracer concentration in cyclonic eddies, consistent with fall 2010 observations of greater *Trichodesmium* abundance in cyclonic eddies. Convergence in cyclones can concentrate buoyant particles such as *Trichodesmium* colonies over periods of several months. The action is slow and persistent. Thus, the model is permissive of the development of different abundance patterns brought about by faster acting mechanisms



over shorter time scales, such as the association with anticyclonic eddies discussed in the next section.

The estimated accumulation rate due to eddy-wind convergence in cyclonic eddies was specific to the conditions simulated and was likely different from the actual accumulation rate. One source of uncertainty is the colony floating velocity, which has been observed to have large variability in the natural environment. We investigated the sensitivity of the results to 10-fold increase and decrease in tracer floating velocity. The accumulation rate increased markedly in the run with increased floating velocity but was not nearly as sensitive to a decrease in floating velocity. This suggests that the factor controlling the accumulation rate is the vertical location of the tracer relative to the sheared velocity structure. An increase in floating velocity skews the vertical tracer distribution higher in the water column (Figure 6), where inward velocities are strongest. A decrease has only a small effect on the vertical tracer distribution and therefore on the rate and direction of transport.

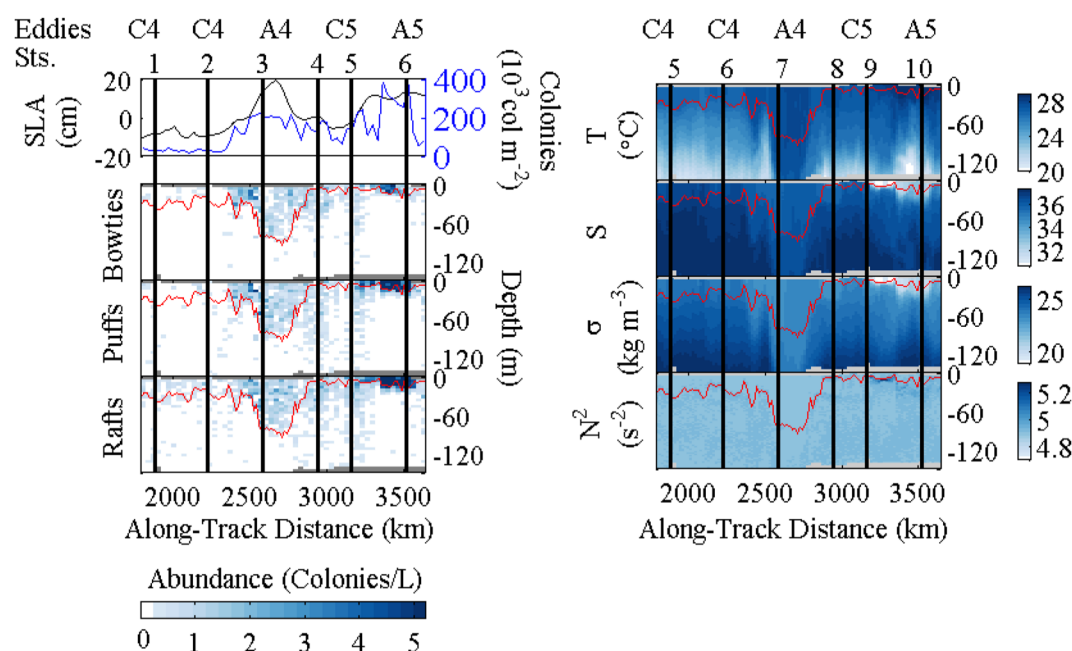
Physical factors that may influence accumulation rates associated with eddy-wind interaction include eddy velocities, wind speed, and diurnal heating/cooling. We initialized the model runs with the physical characteristics of the eddy after the pattern in *Trichodesmium* abundance was already present, which would correspond to the end of the simulation. Based on the eddy decay in the simulation, at an earlier time the observed eddies may have had stronger sea level anomalies and azimuthal velocities. These larger velocities would be associated with greater convergence and colony accumulation rate. A model run at greater eddy strength led to a roughly proportional increase in accumulation rate, consistent with the idealized model discussed in section 3.2. The factor of 1.6 increase produced levels of concentration consistent with the our moderate observation-based estimates.

Quantitatively, there was overlap between the concentration levels reached during the 6 month simulations and those estimated based on observations of the region near Eddy C3 in fall 2010. However, the upper range of the observation-based estimates was higher than simulated levels and unlikely to be achieved by the modeled process alone. Qualitatively, the increase in tracer concentration with proximity to the eddy center is not as smooth in the observations (Figure 2) as in the analytical model (Figure 4) or in the mean section from numerical model (Figure 5). Some differences between our idealized model and observations are to be expected given the simplified nature of the model, which is not intended to reproduce the full complexity of the natural environment.

In addition to wind-driven concentration, spatial heterogeneity in population growth and/or mortality likely contributed to the observed patterns. Many factors can affect net *Trichodesmium* growth rate, such as temperature, salinity, and nutrient availability, as well as mortality. Although this study shows that the physical mechanism proposed herein could contribute a significant portion of the observed enhancement in cyclones, it does not rule out other mechanisms, or interactions between the physical mechanism we have described and biological processes. Furthermore, the time scale over which the physical accumulation occurs is long compared to the doubling time of *Trichodesmium* (2.6–5 days) [Rodier and Le Borgne, 2010; Carpenter *et al.*, 1993]. Thus patterns produced through variable net production of colonies would evolve more quickly than the action of eddy-wind interaction mechanism. Conversely, such biological events may also be more ephemeral in comparison to a physical process that could remain active over the eddy's lifespan.

Another difference between this idealized numerical experiment and the real ocean is in the direction of the wind. In these simulations, the wind vector rotates in direction at a constant rate. While the mechanism we have described is independent of wind direction, the more persistent northeasterly direction of the real-world trade winds in this region could affect the vertical structure of the eddies. Additionally, realistic modulation of wind strength would lead to time variability in the velocities not represented in these idealized simulations. A further dynamic absent from the present simulations is vertical migration, which has been observed in some *Trichodesmium* populations [Villareal and Carpenter, 2003] and investigated with models [White *et al.*, 2006]. The implementation of such a behavioral model in the context of accumulation through eddy-wind interaction could be an interesting avenue of future investigation.

An aspect of the observations not fully explained by the wind-driven concentration model was that although puff colonies were more deeply distributed, their abundance was also elevated in cyclones relative to anticyclones on the fall cruise (Figure 2). Wind-driven concentration would not act to concentrate puffs below the Ekman layer. One possibility is that the puff colonies may have been more buoyant at the time



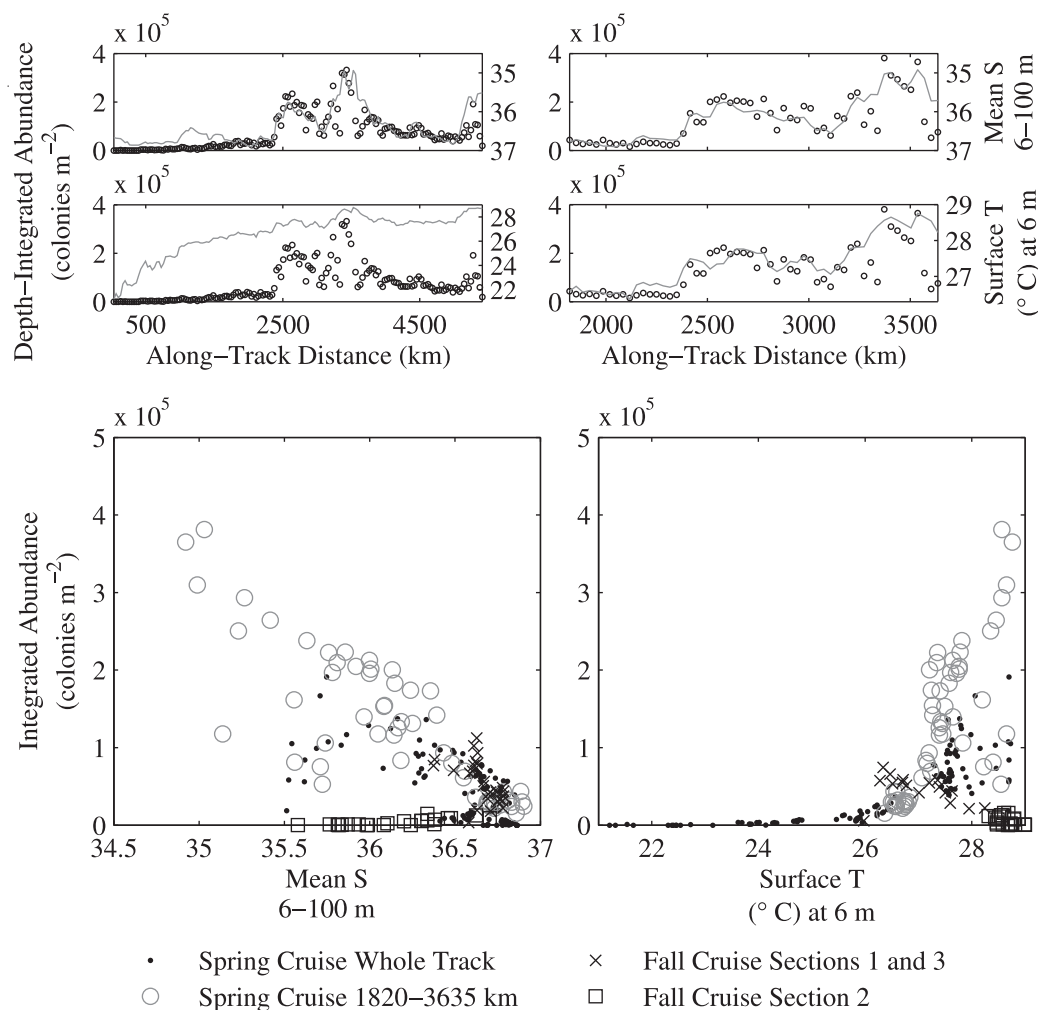
**Figure 8.** *Trichodesmium* abundance and physical properties associated with eddies transected on spring cruise OC471. As in Figure 2, the black line in top plot is sea level anomaly (cm). Blue line in top plot represents integrated raft colony abundance ( $10^3$  colonies  $m^{-2}$ ). Red lines show mixed layer depth as determined by a threshold criterion of  $0.125 \text{ kg m}^{-3}$  density difference from the surface. Right plots (top to bottom) display temperature, salinity, density, and buoyancy frequency.

they were advected toward the center of the cyclone and then began to sink. Villareal and Carpenter [2003] suggest that *Trichodesmium* colonies may employ carbohydrate ballasting to mine phosphorus at the nutricline. Thus, if the colonies became nutrient-stressed when trapped in the convergent cyclone center, the phosphorus mining mechanism could be triggered. The domed isopycnals in cyclones could facilitate this phosphorus acquisition strategy since the phosphocline might be uplifted at the center of the eddy. It is unclear whether puffs and rafts would be affected by nutrient stress at different rates, but rafts are generally observed to be more buoyant and located higher in the water column [Post *et al.*, 2002; Walsby, 1978]. Another mechanism that could explain this pattern is exchange between the puff and raft morphology pools. This could occur if raft-shaped colonies were able to transition to a puff shape or if reproduction of *Trichodesmium* initially in raft-shaped colonies could result in some puff-shaped colonies. In this scenario, raft colonies would converge in the eddy center near the surface, and transition to or produce puff colonies, which would sink to the observed depths. The details of colony morphology determination and transition are not clear and warrant further investigation.

#### 4. *Trichodesmium* Colonies in Fresh Anticyclones: Observations and Discussion

Two anticyclones, A4 (Station 7, Figure 1b) and A5 (Station 10, Figures 1b and 1c), observed in spring 2011, coincided geographically with local peaks in *Trichodesmium* colony abundance (Figure 8). The local correlation between sea level anomaly and integrated *Trichodesmium* abundance was 0.57 ( $p < 0.01$ ) in the region designated Section 2 on the spring cruise, extending from  $18.3^\circ \text{ N}$ ,  $-61.6^\circ \text{ W}$  to  $13.0^\circ \text{ N}$ ,  $-55.6^\circ \text{ W}$  (Figures 1b and 8). Whereas Davis and McGillicuddy [2006] observed elevated colony abundances in warm and salty anticyclones, salinity within these anticyclones was anomalously low. Eddy A5 was strongly stratified, with a fresh lens at the surface. Eddy A4 exhibited a deep mixed layer, with anomalously fresh water evident as deep as 120 m. *Trichodesmium* colonies were also mixed anomalously deep in Eddy A4, down to the mixed layer depth of approximately 80 m (Figure 8); their natural buoyancy likely explains why they were not mixed throughout the full vertical extent of the fresh water anomaly. Locally, abundance patterns were vertically and laterally consistent across morphologies.

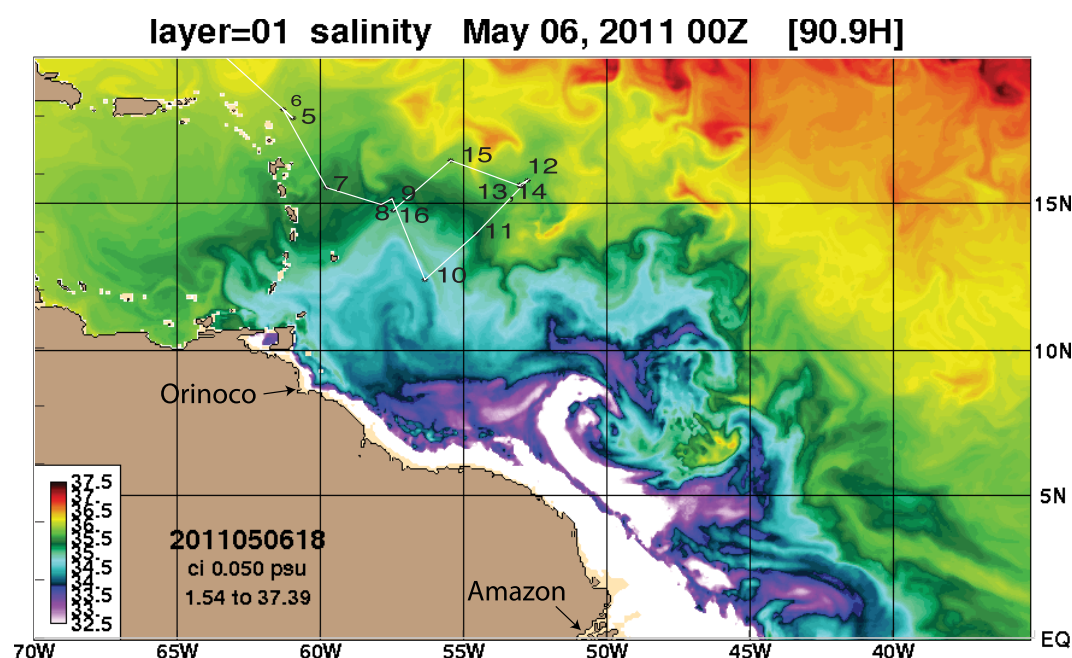
Within spring Section 2, integrated *Trichodesmium* colony abundances were anticorrelated with upper water column salinity and positively correlated with surface temperature. The salinity correlation was strongest



**Figure 9.** Relationship between depth-integrated *Trichodesmium* colony abundance and temperature and salinity. Salinity here is the mean over the upper 100 m, whereas temperature is the near-surface value observed at 6 m depth. (top four plots) (left) temperature, salinity, and colony abundance for the entire spring cruise and (right) for the subset of the spring cruise designated Section 2. Solid line represents (top) salinity or (middle) temperature) and open circles represent *Trichodesmium* colony abundance. Note that salinity axes are reversed in the upper panels to emphasize the degree of anticorrelation between mean salinity and depth-integrated *Trichodesmium* abundance. (bottom) Scatter plots of abundance versus (left) mean salinity and (right) surface temperature. Data from the fall and spring cruises are included in the scatter plots.

when based on average salinity over the upper water column (6–100 m; VPR data were available only below 5 m), which took into account the deeply mixed fresh anomaly in Eddy A4 (correlation coefficient =  $-0.84$ ,  $p < 0.01$ ; Figure 9). In contrast, surface temperature (at 6 m), rather than a depth-mean value, correlated strongly with *Trichodesmium* abundance (correlation coefficient =  $0.76$ ,  $p < 0.01$ ; Figure 9). Rouco *et al.* [2014] found similar correlations using a qPCR method to assess *Trichodesmium* abundance, indicating that the relationships are robust with respect to methodological differences (the VPR samples colonies only, whereas qPCR includes both colonies and free trichomes).

Upper water column salinity and near-surface temperature were themselves correlated over Section 2, but less strongly correlated over the cruise as a whole. On the scale of the full spring cruise, integrated *Trichodesmium* abundances reflect patterns in upper water column salinity more closely than temperature (Figure 9). For the entire cruise, the coefficient of correlation of integrated *Trichodesmium* colony abundance with mean upper 100 m salinity was  $-0.80$  ( $p < 0.01$ ), whereas the correlation with 6 m temperature dropped to  $0.61$  ( $p < 0.01$ ). A similar inverse relationship was described by Borstad [1982] between *Trichodesmium* abundance and salinity at 5 m depth. The relationship between abundance and salinity identified in the spring 2011 data was also present in much of the fall 2010 data. The only exception was fall cruise Section 2, which



**Figure 10.** Sea Surface Salinity from Naval Research Laboratory 1/12° Global HYCOM Hindcast ([http://www7320.nrlssc.navy.mil/GLBHycom1-12/navo/nbrzlsss/nowcast/sss2011050618\\_2011050600\\_909\\_nbrzlsss.001.gif](http://www7320.nrlssc.navy.mil/GLBHycom1-12/navo/nbrzlsss/nowcast/sss2011050618_2011050600_909_nbrzlsss.001.gif)). Cruise track, station locations, and Orinoco and Amazon River mouths are indicated.

exhibited low abundances and a strong association with cyclonic eddies (Figure 2). This subset of the data populated a distinct region in abundance-salinity space (Figure 9).

Previous studies of Amazon plume dispersal point to advection of fresh riverine water into this region along the path of the North Brazil Current (NBC) and Guyana Current [Coles *et al.*, 2013; Richardson and Reverdin, 1987]. The NBC retroflection persists throughout much of the year and sheds rings that translate to the northwest near the South American continent and toward the Caribbean Sea [Fratantoni and Richardson, 2006]. Field [2005] described three examples of rings in which filaments of low salinity Amazon plume water wrapped around the core, and a fourth in which Amazon plume water completely covered the ring. Other studies have described low salinity pools trapped within NBC rings [Hellweger and Gordon, 2002; Kelly *et al.*, 2000]. In the absence of the retroflection, Amazon River water is transported along this pathway directly through the continuation of the NBC in the Guyana Current. In fact, Borstad [1982] linked periodic peaks in *Trichodesmium* abundance to the passage of fresh “pools” or eddies originating from Amazon River Plume water carried by the Guyana Current. Based on  $^{228}\text{Ra}/^{226}\text{Ra}$  ratios sampled in 1981, Moore *et al.* [1986] estimated that 20–30% of surface water near Barbados in May came from the Amazon River estuary.

Plume advection and anticyclonic eddy formation in the NBC region were evident in operational model hindcasts [Chassignet *et al.*, 2009] from May 2011, when spring cruise Eddies A4 and A5 were observed; in fact, in the hindcast for the date anticyclone A5 was transected, a large, fresh anticyclone was present at the same location (Figure 10). Based on objectively analyzed sea level anomaly fields for the period prior to sampling, anticyclone A5 was one of several rings to have been shed into this region from the NBC in spring 2011. Anticyclone A4 may have evolved from remnants of a preceding ring. Its location and history were consistent with a tendency for NBC rings to deflect northward upon interaction with the Lesser Antilles reported by Fratantoni and Richardson [2006]. Such a route might provide clues to the anomalously deep mixing of low salinity water within this eddy, and low salinity pools in the vicinity of Barbados have previously been attributed to remnants of NBC rings [Kelly *et al.*, 2000]. Together, these lines of evidence suggested that the anomalously fresh water encountered in A5 and possibly A4 originated in the Amazon River plume.

Amazon river outflow has been previously investigated as a source of nutrients to the Atlantic. Elevated dissolved iron concentrations have been observed in the Amazon River plume [Bergquist and Boyle, 2006; Tovar-Sanchez and Sañudo-Wilhelmy, 2011; Subramaniam *et al.*, 2008] as well as in a fresh eddy with chemical characteristics typical of Amazon plume water [Fitzsimmons *et al.*, 2013]. Elevated phosphorus has

likewise been observed in the plume [Tovar-Sanchez and Sañudo-Wilhelmy, 2011; Sohm and Capone, 2010; Subramaniam et al., 2008], particularly the SRP and particular organic phosphorus fractions. Demaster and Pope [1996] calculated an intercept of approximately 0.14  $\mu\text{M}$  phosphate at zero nitrate in Amazon shelf waters. At mixing ratios consistent with observed salinities and phosphorus content of 1.10 nmol P colony<sup>-1</sup> [Sañudo-Wilhelmy et al., 2001], this represents more than enough phosphorus to support observed *Trichodesmium* colony abundances.

Phosphorus and iron, the two nutrients believed to be limiting to *Trichodesmium* growth in the tropical and subtropical Atlantic, have thus been demonstrated to be supplied though freshwater of riverine origin, and Amazon outflow has been suggested as a nutrient source to *Trichodesmium* [Coles et al., 2004]. We therefore investigated the relationship between salinity and phosphorus in the spring 2011 data to determine if there was evidence of nutrient supply associated with the fresh water that might explain the observed patterns in *Trichodesmium* abundance (Figure 11). Iron concentrations were not measured in the present study.

Integrated *Trichodesmium* colony abundance from microscopic enumeration at stations from the spring cruise (presented in terms of approximate phosphorus content in Figure 11a) decreased with increasing mean upper water column salinity (correlation coefficient = 0.94,  $p < 0.01$ ), confirming the VPR results. Phosphate measurements from the same stations at which abundance estimates were made (Figure 11b) did not reveal a point-by-point correlation between phosphate concentration and salinity in the upper water column. Rouco et al. [2014] did not find a correlation between phosphate and salinity either, although their analysis was based on a different subset of stations.

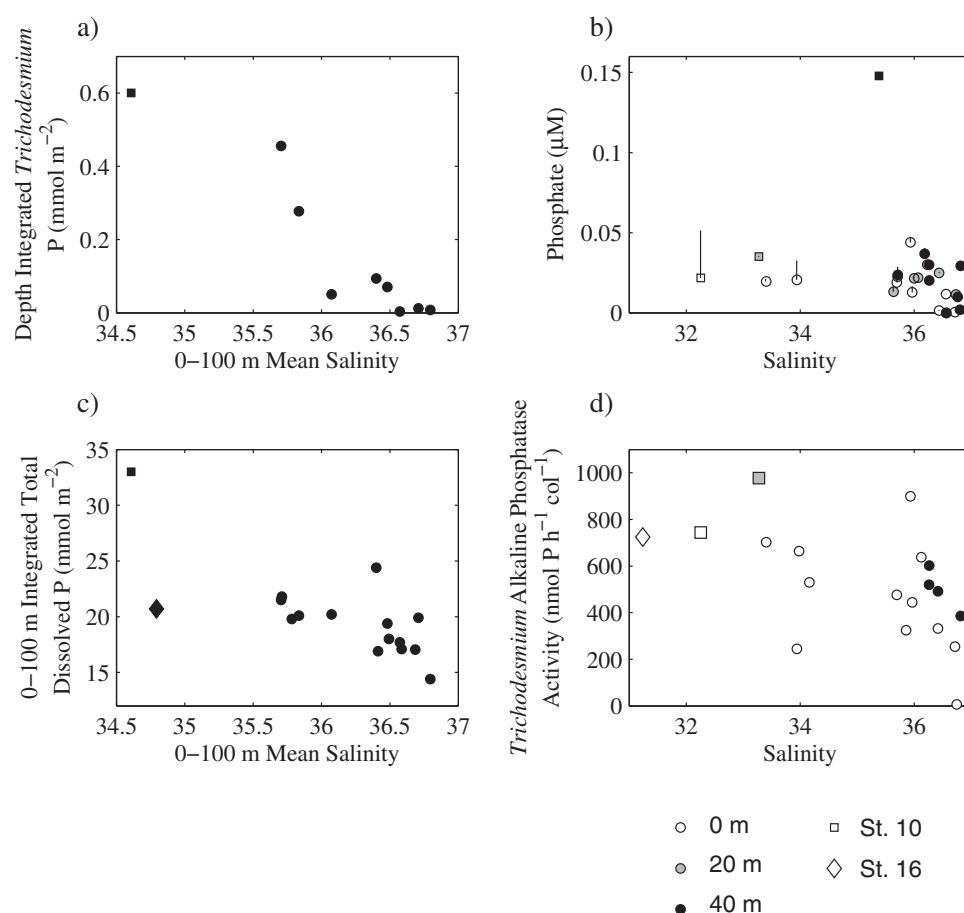
Phosphate levels at low salinities ( $<35$ ) were typically at or slightly above the mean for each sampling depth in the upper 40 m. The exception was Station 16 (not shown), which had very low surface salinity and surface phosphate below the limit of detection. This station was located at the very edge of the fresh-water lens and therefore may represent a distinct ecological and chemical environment. At some low salinity stations, the amount of phosphorus estimated to be present in the form of *Trichodesmium* colonies was comparable to the observed phosphate concentrations (Figure 11b), whereas the ratio of *Trichodesmium* phosphorus to environmental phosphate at higher salinities was lower. This is consistent with a scenario in which the fresh water was a source of phosphorus to *Trichodesmium*, but the phosphorus concentrations have already been drawn down and incorporated into biomass.

*Trichodesmium* alkaline phosphatase activity (APA) was elevated in low salinity waters, (Figure 11d), although the correlation was not significant at the 95% confidence level (correlation coefficient = -0.30,  $p = 0.12$ ). Stations with elevated APA spanned the range of observed surface phosphate concentrations. Although the *Trichodesmium* alkaline phosphatase gene is regulated by P [Orchard et al., 2009], the enzyme activity has not been strongly correlated to phosphate concentration in this region [Orchard et al., 2010]. This is likely in part due to the APA associated with epibionts [Van Mooy et al., 2012]. Regardless, the presence of elevated activity associated with the low salinity waters is permissive of dissolved organic phosphorus (DOP) hydrolysis [Dyhrman et al., 2006; Orchard et al., 2010], which may be rapidly assimilated.

Although the data did not demonstrate a strong relationship between phosphorus and salinity at locations in the upper water column where *Trichodesmium* colonies were observed, 100 m depth-integrated total dissolved phosphorus (TDP) concentrations were higher at lower mean salinities (Figure 11c). In part, the tendency for elevated 100 m integrated TDP at lower salinity may be explained by a shallowing of the phosphocline to the South, in the same region mean 100 m salinity decreased due to the surface lens. However, the relationship is stronger in TDP than phosphate (not shown), highlighting the importance of the organic pool. Taken together, these data suggest that DOP may be an important nutrient source to *Trichodesmium* populations in low salinity waters, exerting influence over their distribution.

Subramaniam et al. [2008] and Goes et al. [2014] described a succession of phytoplankton functional groups downstream in the Amazon River plume. Coastal diatom species dominated until combined N was drawn down, followed by diatom diazotroph associations until Si and soluble reactive phosphorus were drawn down, with *Trichodesmium* further downstream. Stukel et al. [2014] demonstrated the potential contribution of differential grazing along salinity gradients to the pattern of succession as well as the influence of duration of physical retention within an optimal habitat niche on the likelihood of diatom diazotroph association bloom. A key difference between the present study and findings reported by Foster et al. [2007], Goes et al. [2014], and Subramaniam et al. [2008] is that they found higher *Trichodesmium* abundances at "oceanic" (SSS





**Figure 11.** (a) Depth-integrated *Trichodesmium* P content versus upper 100 m mean salinity from spring 2010. P content was estimated from microscopically enumerated colony abundance with conversion factor of 1.10 nmol P per colony [Sañudo-Wilhelmy *et al.*, 2001]. Depth-mean salinity was calculated by trapezoidal integration of measurements at the 20 m intervals (the resolution of the nutrient data). (b) Phosphate concentration versus salinity at matched depths for stations with *Trichodesmium* colony abundance data from microscopic enumeration. Lines extending upward from data points represent the estimated P content of *Trichodesmium* colonies at that depth. (c) Total dissolved phosphorus content, by trapezoidal integration from the surface to 100 m, versus mean salinity based on the same 20 m sampling resolution. (d) *Trichodesmium* alkaline phosphatase activity versus salinity. Samples matched to 0 m salinity were actually taken from net tows at 5 m depth. For Figures 11b and 11d, symbol shading indicates sample depth and symbol shape is used to identify data from two stations, 10 and 16, referred to in the text.

>35) stations compared to “mesohaline” (30 <SSS <35) stations, with most of the elevated nitrogen fixation signal at “mesohaline” stations attributed to diatom-diazotroph associations. In contrast, the greatest *Trichodesmium* abundances in the present study were encountered at stations that would be classified as “mesohaline,” with a significant ( $p < 0.01$ ) correlation between *Trichodesmium* abundance and salinity. The observations discussed by Foster *et al.* [2007], which were also included in the Subramaniam *et al.* [2008] analysis, were made over an area extending from the southern edge of the region sampled in the present study nearly to approximately 5° N, thus closer to the Amazon delta. In addition to this regional difference, although the observations were made at approximately the same time of the year, the Foster *et al.* [2007] observations coincided with a bloom of the *Hemiaulus hauckii*-*Richelia* diatom diazotroph association. That elevated *Trichodesmium* abundances were associated with fresh water influence in the present study suggests that advected river plume water may continue to supply phosphorus and possibly other nutrients necessary for *Trichodesmium* growth after diatom-diazotroph association net growth rates decrease due to reduced Si availability, either through drawdown [Subramaniam *et al.*, 2008; Goes *et al.*, 2014] or dilution by mixing with oceanic water [Stukel *et al.*, 2014].

The supply of phosphorus and iron through river outflow has led to the hypothesis that *Trichodesmium* growth fed by riverine nutrients and growth-conducive surface stratification might seed populations throughout the region. In this context, the present study provides a novel high resolution record of gradients in

*Trichodesmium* colony abundance associated with salinity and indications of eddy-associated transport of *Trichodesmium* biomass and nutrients from the Amazon plume region to the southwestern North Atlantic.

## 5. Conclusions

Two different mesoscale associations between *Trichodesmium* colony abundance and the physical environment were observed in the western tropical North Atlantic. Elevated *Trichodesmium* abundances were attributed to riverine freshwater supply, enhanced within anticyclonic eddies and consistent with the ecological succession described by *Subramaniam et al.* [2008] within river plume waters. Greater colony abundances observed in cyclones were hypothesized to result from gradual but persistent concentration through convergent Ekman transport. The potential action of this mechanism was supported by an analytical model and idealized simulations.

The two opposing mechanisms described might each predominate under different ecological conditions and on different time scales. Under more dynamic bloom conditions, given *Trichodesmium* doubling times in the range of 2.6–5 days [Rodier and Le Borgne, 2010; Carpenter *et al.*, 1993], relationships with factors influencing growth rate, particularly nutrient availability and temperature dependence, are likely to dominate. Enhancement of *Trichodesmium* abundance due to enhanced growth rates in nutrient-rich, fresh, anticyclonic eddies presumably overcomes the more subtle effect of dilution due to eddy-wind driven upwelling and surface divergence. In contrast, the eddy-wind mechanism, while slow, may persist over longer time scales, although for it to predominate, spatial variation in net population growth would have to be small. This requirement may be more likely to be met in uniformly low-nutrient, low-growth conditions, or balanced growth and mortality. Whereas an initial injection of nutrients could quickly lead to an associated signal in abundance, over time, as nutrients were drawn down and growth rates reverted to background levels, the longer-acting physical mechanism might prevail.

Previous studies in other regions have also identified elevated *Trichodesmium* abundances associated with anticyclonic eddies. *Taboada et al.* [2010] attributed their observations to trapping and down gradient transport of *Trichodesmium* biomass in an anticyclonic eddy in the eastern North Atlantic subtropical gyre. *Fong et al.* [2008] and *Davis and McGillicuddy* [2006] attributed enhanced *Trichodesmium* biomass in anticyclonic eddies to nutrient environments conducive to local growth: low nitrogen with sufficient phosphorus and iron. However, the environments and mechanisms producing these conditions were different in each case. The *Fong et al.* [2008] observations were of a decaying anticyclone in the North Pacific subtropical gyre, perhaps associated with upwelling of phosphorus and iron. The *Davis and McGillicuddy* [2006] observations of elevated *Trichodesmium* colony abundance in warm and salty anticyclones in the northern limb of the North Atlantic subtropical gyre have recently been attributed to a combination of down-gradient transport of biomass and phosphorus by anticyclones originating in the Gulf Stream [McGillicuddy, 2014]. The present study attributes enhanced *Trichodesmium* biomass in anticyclonic eddies to phosphorus and iron supply through transport of Amazon plume water to the southwestern North Atlantic. All of these findings underscore the importance of physical processes in shaping abundance patterns underlying the processes driving carbon and nitrogen cycling in the world's tropical and subtropical oceans.

We conclude that wind-driven concentration could produce patterns qualitatively consistent with the association of *Trichodesmium* abundance with cyclones, and that freshwater input was the driving factor in the association with anticyclones. Future work, including analysis of wind-driven concentration through more sophisticated and higher resolution models, could further elucidate the relationship between *Trichodesmium* colonies and their physical environment. In the context of accumulation due to eddy-wind interaction, additional factors to investigate include impacts of more realistic wind forcing, eddy-eddy interaction, and vertical migration. Simulations with a buoyant tracer embedded in a basin-scale eddy-resolving model could facilitate assessment of the mechanism in a more realistic setting.

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